

# Discretising systematically integrable systems

A. RAMANI, *CNRS, Ecole Polytechnique, Palaiseau, France*

in collaboration with

M. MURATA, B. GRAMMATICOS, J. SATSUMA, R. WILLOX

Question asked time and again:

“How do you find a *good* discretisation?”

Standard (rather unsatisfactory) answer:

“With experience and a little bit of luck”

We need a *systematic discretisation approach*

$\exists$  *Infinitely many* discrete analogues of a given continuous system

However for *integrable* systems the answer is almost unique.

Two important names: **Mickens** and **Hirota**

Mickens discretisation rules

- 1 *The orders of “discrete” and “differential” derivatives should be equal*
- 2 *The discrete representations for derivatives must, in general, have nontrivial denominators*
- 3 *Nonlinear terms must be, in general, replaced by nonlocal discrete representations*
- 4 *A property that holds for the differential equation should also be present in the discrete model*

An example: discretise the **Riccati** equation

$$x' = ax^2 + 2bx + f$$

Mickens prescription

$$x' \rightarrow \frac{x_{n+1} - x_n}{\Delta t}$$

$$x^2 \rightarrow x_{n+1}x_n$$

Discrete form

$$x_{n+1} = \frac{(1 + 2b\Delta t)x_n + f\Delta t}{1 - a\Delta tx_n}$$

What about Rule 4?

Integrability by direct linearisation is preserved !

The Hirota method: *bilinearisation* and *gauge invariance*

Riccati example: introduce ansatz

$$x = P/Q$$

Gauge transformation  $P \rightarrow g(t)P$ ,  $Q \rightarrow g(t)Q$  leaves  $x$  invariant

Riccati becomes

$$PQ' - QP' = aP^2 + 2bPQ + fQ^2$$

Gauge-invariance  $\Rightarrow$  nonlocal discretisation of the quadratic terms

$$\frac{Q_{n+1}P_n - P_{n+1}Q_n}{\Delta t} = aP_nP_{n+1} + b(\alpha Q_{n+1}P_n + \beta P_{n+1}Q_n) + fQ_nQ_{n+1}$$

where  $\alpha + \beta = 2$

$$x_{n+1} = \frac{(1 + b\alpha\Delta t)x_n + f\Delta t}{1 - b\beta\Delta t - a\Delta tx_n}$$

Our approach

Discretisation procedure based on

*ad hoc* linearisation of differential system

and Padé-type approximation of the exponential operator

Example, linear first-order equation

$$x' = \alpha x + \beta$$

with solution

$$x(t) = ce^{\alpha t} - \frac{\beta}{\alpha}$$

Time-discretisation

$$x(t + \Delta t) = ce^{\alpha(t+\Delta t)} - \frac{\beta}{\alpha} = e^{\alpha\Delta t} \left( x(t) + \frac{\beta}{\alpha} \right) - \frac{\beta}{\alpha}$$

Rational approximation of exponential

$$e^\sigma = \frac{1 + (\lambda + 1)\sigma}{1 + \lambda\sigma}$$

Finally

$$x_{n+1} = \frac{1 + (\lambda + 1)\alpha\Delta t}{1 + \lambda\alpha\Delta t}x_n + \frac{\Delta t\beta}{1 + \lambda\alpha\Delta t}$$

Second example, Riccati equation

$$x' = (ax + 2b)x + f$$

We find

$$x_{n+1} = \frac{(1 + (\lambda + 1)\Delta t(ax_n + 2b))x_n + f\Delta t}{1 + \lambda\Delta t(ax_n + 2b)}$$

For generic  $\lambda$  not acceptable (violates reversibility)

Taking  $\lambda = -1$  we find

$$x_{n+1} = \frac{x_n + f\Delta t}{1 + 2b\Delta t - a\Delta t x_n} \quad (*)$$

Compare to Hirota result

Equation (\*) is obtained from Hirota for  $\alpha = 0, \beta = 2$

Different derivation

Mapping

$$x_{n+1} = \frac{x_n + f\Delta t}{1 + 2b\Delta t - a\Delta t x_n} \quad (*)$$

can be obtained from

$$x' = (ax + 2b)x + f$$

by ansatz

$$x' \rightarrow (x_{n+1} - x_n)/\Delta t$$

$$x^2 \rightarrow x_{n+1}x_n \quad \text{and} \quad x \rightarrow (x_{n+1} + x_n)/2$$

Two applications ( $\Delta t \equiv \epsilon$ )

The **Lotka-Volterra** system

$$x' = x(\lambda - y) \quad y' = y(x - \mu)$$

Discrete form

$$\frac{x_{n+1}}{x_n} = \frac{1 + (p + 1)(\lambda - y)\epsilon}{1 + p(\lambda - y)\epsilon} \quad \frac{y_{n+1}}{y_n} = \frac{1 + (q + 1)(x - \mu)\epsilon}{1 + q(x - \mu)\epsilon}$$

Freedom on the staggering

$$\frac{x_{n+1}}{x_n} = \frac{1 + (p + 1)(\lambda - y_n)\epsilon}{1 + p(\lambda - y_n)\epsilon} \quad \frac{y_{n+1}}{y_n} = \frac{1 + (q + 1)(x_{n+1} - \mu)\epsilon}{1 + q(x_{n+1} - \mu)\epsilon}$$

Take  $p = -1$  and  $q = 0$

$$\frac{x_{n+1}}{x_n} = \frac{1}{1 - \epsilon\lambda + \epsilon y_n} \quad \frac{y_{n+1}}{y_n} = 1 - \epsilon\mu + \epsilon x_{n+1}$$

**SIR model** for epidemic dynamics

$$S' = -SI \quad I' = -\mu I + SI$$

(Lotka-Volterra with  $\lambda = 0$ )  $\Rightarrow$  discrete form

$$\frac{S_n}{S_{n-1}} = \frac{1 + (p+1)I_n\epsilon}{1 + pI_n\epsilon} \quad \frac{I_{n+1}}{I_n} = \frac{1 + (q+1)(S_n - \mu)\epsilon}{1 + q(S_n - \mu)\epsilon}$$

Staggering different from that of LV

“Intuitive” discretisation

$$\frac{S_n}{S_{n-1}} = \frac{1 + cI_n}{1 + I_n} \quad \frac{I_{n+1}}{I_n} = \frac{a + S_n}{1 + bS_n}$$

Same as “systematic” by  
rescaling of variables and appropriate definition of  $a, b, c$

A system of coupled Riccatis

$$x' = -x^2 + axy$$

$$y' = -y^2 + bxy$$

Painlevé singularity analysis  $\Rightarrow$  5 integrable cases

- i)  $a = 0, b = n$  ( $n$  nonnegative integer)
- ii)  $a = 1, b = 1$
- iii)  $a = 2, b = 2$
- iv)  $a = 1, b = 3$  and its dual  $a = 3, b = 3$
- v)  $a = 1, b = 2$  and its duals  $a = 1, b = 5$  and  $a = 2, b = 5$

Case i) is a special case of the Gambier equation

Discretisation, with  $p = 0$

$$x_{n+1} = \frac{x_n}{1 + x_n}$$

$$y_{n+1} = \frac{(1 + bx_{n+1})y_n}{1 + y_n}$$

Special case of the Gambier mapping:

$$x_{n+1} = \frac{\lambda x_n + \mu}{1 + x_n}$$

$$y_{n+1} = \frac{x_n y_n + \sigma}{1 + \nu y_n}$$

with  $\sigma = 0$ , and specific staggering

For remaining cases again  $p = 0$

$$x_{n+1} = \frac{x_n}{1 + x_n - ay_n}$$

$$y_{n-1} = \frac{y_n}{1 - y_n + bx_n}$$

Study integrability with:

singularity confinement & algebraic entropy

Integrable cases found: *exactly* cases (ii) to (v)

The same values  $a$  and  $b$  lead to integrable for continuous and discrete

Apply our method to the discretisation of **Painlevé equations**

e.g. Painlevé I

$$x'' = x^2 + t$$

Discrete form

$$x_{n+1} + x_n + x_{n-1} = \frac{\alpha n + \beta}{x_n} + 1$$

Was known for 70 years

but only recognised in the 90s

Extend our method to 2nd-order systems

Alas! Not very useful beyond  $P_I$

The Okamoto Hamiltonian formalism for the Painlevé equations  
Hamiltonian is related to the  $\tau$ -function

$$H = (\log \tau)'$$

Write Painlevé equations as Hamiltonian system  
Starting with  $H(x, p, z)$  and equations of motion

$$f(t) \frac{dx}{dt} = \frac{\partial H}{\partial p}$$

$$f(t) \frac{dp}{dt} = -\frac{\partial H}{\partial x}$$

Eliminating  $p$  find equation for  $x$  (and vice versa)  
Miura transformation

Use Hamiltonian formalism for integrable discretisations

Hamiltonian equations of motion are in general of Riccati type

Ansatz for  $x$ :

$$x' \rightarrow x_{n+1} - x_n \quad x^2 \rightarrow x_{n+1}x_n \quad x \rightarrow (x_{n+1} + x_n)/2$$

For  $p$ , analogous ansatz but down-shifted

$$p' \rightarrow p_n - p_{n-1} \quad p^2 \rightarrow p_n p_{n-1} \quad p \rightarrow (p_n + p_{n-1})/2$$

staggering is essential

Discretisation of Painlevé II

Hamiltonian:

$$H(x, p) = \frac{1}{2}p^2 - p \left( x^2 + \frac{t}{2} \right) - \left( \mu + \frac{1}{2} \right) x$$

The equations of motion have the form

$$x' = -x^2 + p - \frac{t}{2} \quad p' = 2xp + \mu + \frac{1}{2}$$

Eliminating  $p$  gives  $P_{II}$  for  $x$

Use ansatz

$$x_{n+1} + x_{n-1} = \frac{x_n(t+2) + \mu + 1/2}{1 - x_n^2}$$

Discrete (autonomous)

Deautonomisation: here take  $t$  linear in  $n$

## Painlevé III

$$H(x, p) = 2x^2p^2 - p(zx^2 + 2\mu x - z) + \kappa zx$$

with  $z = e^t$  and  $f(z) = 1$

$$x' = x^2(4p - z) - 2\mu x + z$$

$$p' = -4p^2x + 2p(xz + \mu) - \kappa z$$

## Discretisation

$$x_{n+1}x_{n-1} = \frac{x^2(\mu^2 - 1) - 2\mu xz + z^2}{x^2z^2 + 2xz(\mu - 2\kappa) + \mu^2 - 1}$$

with  $z = \lambda^n$  we find  $q$ -discrete  $P_{\text{III}}$

Painlevé IV

$$H(x, p) = 2xp^2 - p(x^2 + 2tx + \mu) + \kappa x$$

and

$$x' = -x^2 + 2x(2p - t) - \mu$$

$$p' = -2p^2 + 2p(x + t) - \kappa$$

Discretisation

$$(x_{n+1} + x_n)(x_n + x_{n-1}) = \frac{(x^2 - \mu)^2 - 4x^2}{(x + t)^2 - 2\kappa - 1}$$

Deautonomisation  $\Rightarrow t$  linear in  $n$

Painlevé V

$$H(x, p) = x(x - 1)^2 p^2 - p(\nu(x - 1)^2 - \mu x(x - 1) - zx) + \kappa(x - 1)$$

and  $z = e^t$

$$x' = 2px^3 - (4p + \nu - \mu)x^2 + (2p + z + 2\nu - \mu)x - \nu$$

$$p' = -p^2(3x^2 - 4x + 1) + p(2(\nu - \mu)x + \mu - 2\nu - z) - \kappa$$

Equation for  $x$  is *not* of Riccati type

Introduce auxiliary variable  $u = xp$  and eliminate  $p$

$$x' = 2ux^2 - 4ux - (\nu - \mu)x^2 + 2u + (z + 2\nu - \mu)x - \nu$$

$$u' = - \left( x - \frac{1}{x} \right) u^2 + (\nu - \mu)xu - \frac{\nu}{x}u - \kappa x$$

Discretisation  $\Rightarrow$  discrete equation for  $x$

but *not* in canonical form

Introduce new variable

$$x_n = \frac{y_n - 1}{y_n + 1}$$

Mapping for  $y$

$$\begin{aligned} & (y_{n+1}y_n - 1)(y_n y_{n-1} - 1) \\ &= \frac{(y_n^2 - 1)^2(\mu^2 - 4) + 4y_n(y_n - 1)^2(4\kappa + 2\mu\nu - \mu^2) + 16\nu^2 y_n^2}{(zy_n - \mu)^2 - 4} \end{aligned}$$

with  $z = \lambda^n$  we find  $q$ -discrete  $P_V$

## Painlevé VI

$$H(x, p) = x(x-1)(x-t)p^2 - p(\nu(x-1)(x-t) + \rho x(x-t) + \mu x(x-1)) + \kappa(x-t)$$

where for the time being we do not care about  $f(t)$

$$x' = 2px^3 - (2p(t+1) + \nu + \mu + \rho)x^2 + (2pt + \nu(t+1) + \mu + \rho t)x - \nu t$$

$$p' = -p^2(3x^2 - 2x(t+1) + t) + p(2(\nu + \mu + \rho)x - \mu - \nu(t+1) - \rho t) - \kappa$$

Again introduce  $u = xp$  and eliminate  $p$

$$x' = 2ux^2 - (\nu + \mu + \rho)x^2 - 2(t+1)ux + (\nu(t+1) + \mu + \rho t)x + 2ut - \nu t$$

$$u' = - \left( x - \frac{t}{x} \right) u^2 + (\nu + \mu + \rho)xu - \frac{\nu t}{x}u - \kappa x$$

The mapping for  $x$  is not in canonical form

$$x_n = \sqrt{t} \frac{1 - y_n}{1 + y_n}$$

The (continuous) independent variable must also be changed

$$t = \left( \frac{1 - s}{1 + s} \right)^2$$

We finally find ( $s = \lambda^n$ ,  $\sigma = \rho + \mu$ )

$$\begin{aligned} & \frac{(y_{n+1}y_n - s^2)(y_n y_{n-1} - s^2)}{(y_{n+1}y_n - 1)(y_n y_{n-1} - 1)} \\ &= \frac{(\rho(y_n - s)^2 + \mu(y_n + s)^2)^2 - 4t^{-1}(y_n^2 - s^2)^2}{(\sigma^2 - 4t^{-1})(y_n^2 - 1)^2 + (16\kappa - 4\sigma(\sigma + 2\nu))y_n(y_n - 1)^2 + 16\nu^2 y_n^2} \end{aligned}$$

All Painlevé equations could be discretised

*An intriguing remark:*

only the discrete forms of the “standard” family were obtained

Why? The standard forms are not even the more fundamental!

On the contrary, if we implement full freedom

$P_I \rightarrow P_{II}$ ,  $P_{II} \rightarrow P_{III}$ ,  $P_{III} \rightarrow P_{VI}$ ,  $P_{IV} \rightarrow P_{VI}$  and  $P_V, P_{VI} \rightarrow$  higher

More important

Where are the other discrete forms of the Painlevé equations?

We must find a different approach

A (not so) short introduction to the QRT mapping

Motivation:

Autonomous limit of Painlevé transcendents  $\Rightarrow$  elliptic functions

Angle of attack:

*To obtain discrete Painlevé equations*

start from mapping with elliptic function solutions

then extend by deautonomisation

Enter QRT mapping

Our strategy for discretisation:

*Perform discretisation on autonomous form*

*require integrability i.e. ask that they be of QRT type*

*if non-autonomous form is already known, identify it*

*if not, deautonomise maintaining integrability*

Ansatz for  $x$ :

$$x'' \rightarrow x_{n+1} + x_{n-1} - 2x_n \quad x \rightarrow a_1(x_{n+1} + x_n) + a_2x_n$$

$$x^2 \rightarrow b_1x_{n+1}x_{n-1} + b_2x_n(x_{n+1} + x_{n-1}) + b_3x_n^2$$

## Canonical forms of $A_1$ QRT matrices

$$(I) \quad A_1 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} \quad x_{n+1} + x_{n-1} = F(x_n)$$

$$(II) \quad A_1 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad x_{n+1}x_{n-1} = F(x_n)$$

$$(III) \quad A_1 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \quad (x_{n+1} + x_n)(x_n + x_{n-1}) = F(x_n)$$

$$(IV) \quad A_1 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -1 \end{pmatrix} \quad (x_{n+1}x_n - 1)(x_nx_{n-1} - 1) = F(x_n)$$

$$(V) \quad A_1 = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 2 & z \\ 1 & z & 0 \end{pmatrix} \quad \frac{(x_{n+1} + x_n + z)(x_n + x_{n-1} + z)}{(x_{n+1} + x_n)(x_n + x_{n-1})} = F(x_n)$$

$$(VI) \quad A_1 = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -z - 1 & 0 \\ 0 & 0 & z \end{pmatrix} \quad \frac{(x_{n+1}x_n - z)(x_nx_{n-1} - z)}{(x_{n+1}x_n - 1)(x_nx_{n-1} - 1)} = F(x_n)$$

$$(VII) \quad A_1 = \begin{pmatrix} 0 & 0 & 1 \\ 0 & -2 & -2z \\ 1 & -2z & z^2 \end{pmatrix} \quad \frac{(x_{n+1} - x_n - z)(x_{n-1} - x_n - z) + 4zx_n}{x_{n+1} - 2x_n + x_{n-1} - 2z} = F(x_n)$$

(VIII)

$$A_1 = \begin{pmatrix} 0 & 0 & 1 \\ 0 & z + 1/z & 0 \\ 1 & 0 & 1 \end{pmatrix} \frac{(x_{n+1}z + x_n)(x_{n-1}z + x_n) - z^2}{(x_{n+1} + zx_n)(x_{n-1} + zx_n) - 1} = F(x_n)$$

Matrix  $A_1$  for case VI has different structure with  $\alpha_1 \neq 0$

Traditionally associated with the Painlevé VI equation

but it is possible to transform to  $\alpha_1 = 0$

(VI')

$$A_1 = \begin{pmatrix} 0 & 0 & 1 \\ 0 & z + 1/z & 0 \\ 1 & 0 & 0 \end{pmatrix} \frac{(x_{n+1} + zx_n)(zx_n + x_{n-1})}{(zx_{n+1} + x_n)(x_n + zx_{n-1})} = F(x_n)$$

More convenient for degeneration process starting from VIII

## Discretising the Painlevé I equation

$$x'' = x^2 + \lambda x + \mu$$

Ansatz

$$ax_{n+1}x_{n-1} + bx_n(x_{n+1} + x_{n-1}) + cx_n^2 + f(x_{n+1} + x_{n-1}) + gx_n + h = 0$$

*Constraints*

If  $a = 0$ ,  $b = 0$  and  $c = 0$ , mapping becomes linear

If  $a = 0$ ,  $b = 0$  and  $f = 0$ , mapping becomes linear

If  $a = 0$  and  $b^2h - bfg + cf^2 = 0$ , factorisation and mapping becomes linear

If  $ac - b^2 = 0$ ,  $ag - 2bf = 0$  and  $ah - f^2 = 0$  factorisation into linear first-order mappings

A few selected examples

Two difference P<sub>IS</sub>

$$x_{n+1} + x_{n-1} = -\frac{h}{x_n + f}$$

$$x_{n+1} + x_n + x_{n-1} = -g - \frac{h}{x_n}$$

A  $q$ -P<sub>I</sub>

$$x_{n+1}x_{n-1} = -gx_n - h$$

Also a P<sub>I</sub> from the P<sub>IV</sub> family

$$(x_{n+1} + x_n)(x_n + x_{n-1}) = -gx_n - h$$

From the family of  $P_{\text{VI}}$ , two new equations

$$\left( \frac{x_{n+1} + z_n z_{n+1} x_n}{z_n z_{n+1} x_{n+1} + x_n} \right) \left( \frac{z_n z_{n-1} x_n + x_{n-1}}{x_n + z_n z_{n-1} x_{n-1}} \right) = \frac{x_n z_{n-1} z_n^2 z_{n+1} - 1}{x_n - z_{n-1} z_n^2 z_{n+1}}$$

and

$$\left( \frac{x_{n+1} + z_n z_{n+1} x_n}{z_n z_{n+1} x_{n+1} + x_n} \right) \left( \frac{z_n z_{n-1} x_n + x_{n-1}}{x_n + z_n z_{n-1} x_{n-1}} \right) = \frac{1}{z_n} \frac{x_n z_{n-1} z_n z_{n+1} + 1}{-x_n + z_{n-1} z_n z_{n+1}}$$

where  $z_n = z_0 \lambda^n$

For the difference  $E_8^{(1)}$  case

$$\frac{(x_n - x_{n+1} + \zeta_n^2)(x_n - x_{n-1} + \zeta_{n-1}^2) + 4x_n\zeta_n\zeta_{n-1}}{\zeta_{n-1}(x_n - x_{n+1} + \zeta_n^2) + \zeta_n(x_n - x_{n-1} + \zeta_{n-1}^2)} = \zeta_n + \zeta_{n-1} - \frac{x_n + f}{\zeta_n + \zeta_{n-1}}$$

$$\frac{(x_n - x_{n+1} + \zeta_n^2)(x_n - x_{n-1} + \zeta_{n-1}^2) + 4x_n\zeta_n\zeta_{n-1}}{\zeta_{n-1}(x_n - x_{n+1} + \zeta_n^2) + \zeta_n(x_n - x_{n-1} + \zeta_{n-1}^2)} = z_{n+1} + z_n + z_{n-1} - \frac{x_n + f}{z_{n+1} + z_n + z_{n-1}}$$

where  $z_n = z_0 + \lambda n$  and  $\zeta_n = z_n + z_{n+1}$

And for the multiplicative  $E_8^{(1)}$  case

$$\frac{(x_{n+1}\zeta_n + x_n)(x_{n-1}\zeta_{n-1} + x_n) - \phi_n}{(x_{n+1} + \zeta_n x_n)(x_{n-1} + \zeta_{n-1} x_n) - \phi_n / (\zeta_n \zeta_{n-1})} = \frac{x_n - f\zeta_n \zeta_{n-1}}{x_n \zeta_n \zeta_{n-1} - f}$$

$$\frac{(x_{n+1}\zeta_n + x_n)(x_{n-1}\zeta_{n-1} + x_n) - \phi_n}{(x_{n+1} + \zeta_n x_n)(x_{n-1} + \zeta_{n-1} x_n) - \phi_n / (\zeta_n \zeta_{n-1})} = \frac{x_n z_n + g\zeta_n \zeta_{n-1}}{-x_n \zeta_n \zeta_{n-1} / z_n + g}$$

where  $z_n = z_0 \lambda^n$ ,  $\zeta_n = z_n z_{n+1}$  and  $\phi_n = (\zeta_n^2 - 1)(\zeta_{n-1}^2 - 1)$

## Results on PII

A contiguity of continuous  $P_{VI}$

$$\frac{z_{n+1} + z_n}{x_n + x_{n+1}} + \frac{z_n + z_{n-1}}{x_n + x_{n-1}} = \frac{2z_n}{x_n + a} + \frac{2z_n}{x_n + b}$$

A  $E_8^{(1)}$  example

$$\frac{(x_{n+1}\zeta_n + x_n)(x_{n-1}\zeta_{n-1} + x_n) - \phi_n}{(x_{n+1} + \zeta_n x_n)(x_{n-1} + \zeta_{n-1} x_n) - \phi_n / (\zeta_n \zeta_{n-1})} = \frac{x_n^2 + dz_n^2 x_n - z_n^8 - 1 - fz_n^4}{x_n^2 z_n^4 + dz_n^2 x_n - z_n^4 - 1/z_n^4 - f}$$

Many more exist

## Conclusion

The new method is very powerful

It allows a “bottom-up” construction of Painlevé equations

First results of equations described by  $E_8^{(1)}$  Weyl group

Not presented results: **linearisable** equations

## Remaining tasks

Find full freedom of the new equations

Derive discrete forms of the other Painlevé equations