## LECTURES 5 AND 6

## 1. Lecture 5 (First Chern form, Kähler connection, and curvature)

Suppose $\tilde{h}$ is another Hermitian metric. Note that $\frac{\tilde{h}_{\alpha}}{h_{\alpha}}=\frac{\tilde{h}_{\beta}}{h_{\beta}}$. Hence $\tilde{h}=h e^{-\phi}$ where $\phi$ is some smooth globally defined function. Note that $c_{1}(\tilde{h})=c_{1}(h)+\frac{\sqrt{-1}}{2 \pi} \partial \bar{\partial} \phi$. In other words, $\left[c_{1}(h)\right]$ (the De Rham cohomology class) is independent of the metric chosen !! This topological quantity is called the first Chern class of the line bundle $L$.

It seems that the first Chern form almost gives us a Kähler form except for the point that it may not be positive-definite. We define the following : A holomorphic line bundle is said to positive is $c_{1}(L)(h)$ is positive for some Hermitian metric $h$ (likewise, negative). A ( 1,1 ) De Rham cohomology class [ $\omega$ ] is said to be positive if there is a Kähler form in it, i.e., there is a Kähler form $\omega^{\prime}$ such that $\omega^{\prime}=\omega+d \eta$.

Exercise : Using the $\partial \bar{\partial}$ lemma, show that if $[\omega]=\left[c_{1}(L)\right]$ then there is a smooth Hermitian metric $h$ so that $c_{1}(h)=\omega$, i.e., every form in the first Chern class can be realised using a Hermitian metric.

By the way,

Theorem 1.1 (Kodaira's embedding theorem). A compact complex manifold can be holomorphically embedded as a submanifold in $\mathbb{C P}^{n}$ iff it has a positive holomorphic line bundle $L$.

The Fubini-Study metric (up to a factor) as we defined it is the first Chern form of $O(1)$ equipped with a metric. Before we look at that, here are a couple of points :
(1) Recall that if $L$ is a holomorphic bundle, then there is a dual bundle $L^{*}$. It is defined set theoretically as $\cup_{p} L_{p}^{*}$ and the topology and complex structure are given by local trivialisations, i.e., if $s_{\alpha}$ are local holomorphic bases for $L$, then $s_{\alpha}^{*}$ defined as $s_{\alpha}^{*}\left(s_{\alpha}\right)=1$ are local holomorphic bases of $L^{*}$. (So the transition functions are $\tilde{g}_{\alpha \beta}=1 / g_{\alpha \beta}$. (More generally, if $E$ is a holomorphic (or even smooth for that matter) vector bundle, then $E^{*}=\cup_{p} E_{p}^{*}$ set theoretically. The topology and complex structure are given as: If $e_{\alpha, i}$ are local holomorphic bases for $E$ then $e_{\alpha}^{* i}\left(e_{\alpha, j}\right)=\delta_{j}^{i}$ are holomorphic local bases of $E^{*}$. So the transition functions are $\left(\left[g_{\alpha \beta}\right]^{-1}\right)^{T}$.)
(2) If $h$ is a Hermitian metric on $L$, then there is a natural Hermitian metric $h^{*}$ on $L^{*}$ defined as $h_{\alpha}^{*}=1 / h_{\alpha}$. Indeed, $h_{\alpha}^{*}=h_{\beta}^{*}\left|\tilde{g}_{\alpha \beta}\right|^{2}$ and hence $h^{*}$ is a well-defined metric. Therefore, $c_{1}\left(h^{*}, L^{*}\right)=-c_{1}(h, L)$.
There is an obvious Hermitian metric on $O(-1) \subset \mathbb{C P}^{n} \times \mathbb{C}^{n+1}$ coming from the Euclidean metric on $\mathbb{C}^{n+1}$. Locally, in the chart $U_{0}$, if $s_{0} \in O(-1)=\left(1, z^{1}, z^{2}, \ldots\right)$ is a local basis, then $h_{\alpha}=1+|z|^{2}$. Hence, $c_{1}(O(-1), h)=-\frac{\sqrt{-1}}{2 \pi} \partial \bar{\partial} \ln \left(1+|z|^{2}\right)$ and $c_{1}\left(O(1), h^{*}\right)=\omega_{F S}=\frac{\sqrt{-1}}{2 \pi} \partial \bar{\partial} \ln \left(1+|z|^{2}\right)$. The associated Hermitian metric is $h_{i \bar{j}}=\frac{1}{\pi} \frac{\partial^{2}}{\partial z^{i} \partial \bar{z} \bar{j}} \ln \left(1+|z|^{2}\right)$.

Given the Fubini-Study metric, we can calculate the induced metric on submanifolds. Indeed, let $F\left(X^{i}\right)$ be a degree $d$ homogeneous polynomial such that $\nabla F \neq 0$ on $S=F^{-1}(0)$. As we saw earlier, $S$ is a compact complex submanifold of $\mathbb{C P}^{n}$. Assume without loss of generality that $\frac{\partial F}{\partial X^{1}} \neq 0$
near a point on $S$ where $X^{0} \neq 0$. Therefore, locally, $z^{1}=f\left(z^{2}, \ldots, z^{n}\right)$ and $F\left(1, f, z^{2}, \ldots\right)=0$. Now $\omega_{F S} \left\lvert\, S=\frac{\sqrt{-1}}{2 \pi} \partial \bar{\partial} \ln \left(1+|f|^{2}+\sum_{i=2}\left|z^{i}\right|^{2}\right)\right.$, which equals

$$
\begin{gather*}
\frac{1}{1+|f|^{2}+\sum_{i=2}\left|z^{i}\right|^{2}} \frac{\sqrt{-1}}{2 \pi}\left(\partial f \wedge \overline{\partial f}+\sum_{i=2} d z^{i} \wedge d \bar{z}^{i}-\frac{\left.\sum_{i, j}\left(\bar{f} \frac{\partial f}{\partial z^{i}}+\bar{z}^{i}\right)\left(f \frac{\partial \bar{f}}{\partial \bar{z}^{j}}+z^{j}\right) d z^{i} \wedge d \bar{z}^{j}\right)}{1+|f|^{2}+\sum_{i=2}\left|z^{i}\right|^{2}}\right) \\
=\frac{1}{1+|f|^{2}+\sum_{i=2}\left|z^{i}\right|^{2}} \frac{\sqrt{-1}}{2 \pi}\left(\delta_{i \bar{j}}+a_{i} \bar{a}_{j}-c_{i} \bar{c}_{j}\right) d z^{i} \wedge d \bar{z}^{j} \tag{1.1}
\end{gather*}
$$

Noting that $\omega^{n}=\operatorname{det}\left(h_{i j}\right) n!\frac{\sqrt{-1}}{2} d z^{1} \wedge d \bar{z}^{1} \ldots$, we can calculate the volume form by computing the determinant of the above matrix.

Exercise : Calculate the determinant of the matrix above.
Before we go on further, we need another notion from vector bundles. If $E$ is a rank-r vector bundle, then $\Lambda^{r} E$ is also a vector bundle where $\Lambda_{p}^{r} E_{p}=E_{p} \wedge E_{p} \wedge \ldots$ If $e_{\alpha, i}$ is a local basis of $E$, then $\eta_{\alpha}=e_{\alpha, 1} \wedge e_{\alpha, 2} \ldots e_{\alpha, r}$ is a local basis for $\Lambda^{r} E$. If $e_{\beta, i}=\left[g_{\alpha \beta}\right]_{i}^{j} e_{\alpha, j}$, then

Exercise : Prove that $\eta_{\beta}=\operatorname{det}\left(g_{\alpha \beta}\right) \eta_{\alpha}$
Therefore, if $H$ is a Hermitian metric on $E$, then $\operatorname{det}(H)$ is a Hermitian metric on $\operatorname{det}(E)$. In particular, $\operatorname{det}\left(T^{* 1,0} M\right)=K_{M}$ is called the canonical bundle of $M$ and if $\omega$ is a Kähler form on $M$, then $\operatorname{det}(h)^{-1}$ (which is basically $\left.\left(\frac{\omega^{n}}{n!}\right)^{-1}\right)$ is a Hermitian metric on $K_{M}$. Thus, $c_{1}\left(K_{M}, \operatorname{det}(h)\right)=\frac{\sqrt{-1}}{2 \pi} \partial \bar{\partial} \ln (\operatorname{det}(h))$.

Exercise : Compute $c_{1}\left(K_{\mathbb{C P}^{n}}\right)$ with the metric induced from the Fubini-Study metric. A harder exercise is to compute $c_{1}\left(K_{S}\right)$ with the metric above. If you do the calculations correctly, you should get something like $c_{1}\left(K_{\mathbb{C P}}^{n}\right)=-(n+1) \omega_{F S}$, and $c_{1}\left(K_{S}\right)=(d-n-1) \omega_{F S}+\sqrt{-1} \partial \bar{\partial} \psi$ for some smooth globally defined function $\psi$.

The above exercises show that $K_{\mathbb{C P}^{n}}$ is a negative line bundle, and if $d$ is large, then $K_{S}$ is a positive line bundle (and hence $K_{S}^{*}$ is negative). When $d=n+1,\left[c_{1}\left(K_{S}\right)\right]=[0]$. Such an $S$ (for example $\left(X^{0}\right)^{5}+\ldots\left(X^{4}\right)^{5}=0$ called the Fermat quintic) is called a Calabi-Yau manifold. These manifolds are important in String theory. They have nice Kähler metrics with good curvature properties.

## 2. Lecture 6 (Kähler connection and curvature)

Now we shall study the curvature of the Levi-Civita connection of Kähler manifolds. Recall that the Levi-Civita connection $\nabla$ gives us a way to find the directional derivative of vector fields. $\nabla_{X} Y$ is the derivative of $Y$ along $X$. It is uniquely determined by a few properties. More generally, given a smooth vector bundle $E$ on a smooth manifold $M$, a connection is way to find the directional derivative of sections. It is defined as a map :

$$
\nabla: \text { Smooth sections of } E \times \text { Smooth vector fields on } M \rightarrow \text { Smooth sections of } E
$$

satisfying
(1) $\nabla_{f_{1} X_{1}+f_{2} X_{2}} s=f_{1} \nabla_{X_{1}} s+f_{2} \nabla_{X_{2}} s$ where $f_{1}, f_{2}$ are smooth functions on $M$.
(2) $\nabla_{X}\left(s_{1}+s_{2}\right)=\nabla_{X} s_{1}+\nabla_{X} s_{2}$.
(3) $\nabla_{X}(f s)=X(f) s+f \nabla_{X} s$.

If there is a Hermitian metric $H$ on $E$, then $\nabla$ is said to be metric compatible if $X\left(H\left(s_{1}, s_{2}\right)\right)=$ $H\left(\nabla_{X} s_{1}, s_{2}\right)+H\left(s_{1}, \nabla_{X} s_{2}\right)$.

The Levi-Civita connection can be used to differentiate not just vector fields, but also induces a connection on $T^{*} M$ by $X(\omega(Y))=\nabla_{X} \omega(Y)+\omega\left(\nabla_{X} Y\right)$. The Christoffel symbols are defined as $\nabla_{\partial_{k}} \partial_{j}=\Gamma_{j k}^{i} \partial_{i}$. For the Levi-Civita connection there is a nasty formula for these beasts. But it is much simpler for calculations to note that the Christoffel symbols at $p$ vanish in normal coordinates near $p$. For one-forms, $\nabla_{\partial_{k}} d x^{j}\left(\partial_{i}\right)=\partial_{k}\left(\delta_{i}^{j}\right)-\delta_{l}^{j} \Gamma_{k i}^{l}$ and hence $\nabla_{\partial_{k}} d x^{j}=-\Gamma_{k i}^{j} d x^{i}$. Using these two connections, we can talk about differentiating other tensors. For instance if $J=J_{j}^{i} d x^{j} \otimes \partial_{i}$, then $\nabla_{X} J$ is defined to be $\nabla_{X} J=X\left(J_{j}^{i}\right) d x^{j} \otimes \partial_{i}-J_{j}^{i} \Gamma_{k l}^{j} X^{k} d x^{l} \otimes \partial_{i}+J_{j}^{i} d x^{j} \otimes \Gamma_{i l}^{k} X^{l} \partial_{k}$.

The Levi-Civita connection $\nabla$ on a complex manifold can be extended complex linearly to a connection on $\mathbb{C T M}$. On a Kähler manifold, since there are holomorphic normal coordinates at every point $p, \nabla J=0$ (because $J$ has constant coefficients). We can define the Christoffel symbols in the $z, \bar{z}$ basis as follows.

$$
\begin{equation*}
\nabla_{\partial_{z^{k}}} \partial_{z^{j}}=\Gamma_{j k}^{i} \partial_{z^{i}}+\Gamma_{j k}^{i} \partial_{\bar{z}^{i}} \tag{2.1}
\end{equation*}
$$

and so on. Since $J \partial_{z^{i}}=\sqrt{-1} \partial_{z^{i}}$ and $\nabla J=0$, we see that $J \nabla \partial_{z^{i}}=\sqrt{-1} \nabla \partial_{z^{i}}$. Therefore, $\nabla_{\partial_{z^{k}}} \partial_{z^{i}} \in T^{1,0}$ and hence $\Gamma_{j k}^{\bar{i}}=0$. Now the torsion-free ness forces $\nabla_{\bar{i}} \partial_{k}=\nabla_{k} \partial_{\bar{i}}$ and hence both of them vanish. So the only surviving symbols are $\Gamma_{j k}^{i}=\Gamma_{k j}^{i}$ and $\Gamma_{j k}^{i}=\overline{\Gamma_{j k}^{i}}$. One can now calculate the Christoffel symbols. The Levi-Civita connection satisfies $\nabla g=0$. Since $\nabla J=0, \nabla h=0$.

$$
\begin{gather*}
h_{i \bar{j}, k}=\partial_{k}\left(h\left(\partial_{i}, \partial_{\bar{j}}\right)\right)=h\left(\nabla_{k} \partial_{i}, \partial_{\bar{j}}\right)+h\left(\partial_{i}, \nabla_{k} \partial_{\bar{j}}\right) \\
=h\left(\Gamma_{k i}^{l} \partial_{l}, \partial_{\bar{j}}\right)=\Gamma_{k i}^{l} h_{l \bar{j}} \\
\Rightarrow \Gamma_{j k}^{i}=h^{i \bar{i}} \partial_{j} h_{k \bar{l}} . \tag{2.2}
\end{gather*}
$$

The upper indices indicate the inverse. So the matrix of 1 -forms $\Gamma_{k}^{i}=\left[\partial h h^{-1}\right]_{k}^{i}$. The Riemann curvature tensor (extended $\mathbb{C}$-linearly) is

$$
\begin{equation*}
\left(\nabla_{k} \nabla_{\bar{l}}-\nabla_{\bar{l}} \nabla_{k}\right) \partial_{i}=R_{i, k \bar{l}}^{-j--} \partial_{j} . \tag{2.3}
\end{equation*}
$$

Exercise : Prove that the other covariant derivatives commute.
Upon computing

$$
\begin{align*}
& \left(\nabla_{k} \nabla_{\bar{l}}-\nabla_{\bar{l}} \nabla_{k}\right) \partial_{i}=-\nabla_{\bar{I}}{ }^{\Gamma}{ }_{k i}^{j} \partial_{j} \\
& =-\partial_{\bar{l}}\left(\Gamma_{k i}^{j}\right) \partial_{j}=-\partial_{\bar{l}}\left(h^{j} \partial_{k} h_{i \bar{l}}\right) . \tag{2.4}
\end{align*}
$$

In holomorphic normal coordinates, $R_{i, k \bar{l}}^{-j-}(p)=-\partial_{\bar{l}} \partial_{k} h_{i j}(p)$. By the local $\partial \bar{\partial}$ lemma, since $h_{i \bar{j}}=\partial_{\bar{j}} \partial_{i} \phi$ for some smooth $\phi$, we can interchange the derivatives to get many symmetries of the holomorphic Riemann curvature tensor. Recall that the Ricci curvature in usual Riemannian geometry is defined as $\operatorname{Ricc}(Y, Z)=\operatorname{tr}(X \rightarrow R(X, Y) Z)$. The Ricci tensor is symmetric.

In holomorphic normal coordinates, the Ricci tensor can be computed (extended $\mathbb{C}$-linearly) as

$$
\begin{gather*}
\operatorname{Ricc}_{i \bar{l}}=R_{i_{-j} \bar{l}}^{-j_{-}}(p)=-\sum_{j} \partial_{\bar{l}} \partial_{j} h_{i j}(p) \\
=-\sum_{j} \partial_{\bar{l}} \partial_{j} \partial_{i} \partial_{j} \phi(p)=-\sum_{j} \partial_{\bar{j}} \partial_{j} \partial_{i} \partial_{\bar{l}} \phi(p)=-\sum_{j} \partial_{\bar{l}} \partial_{i} h_{j \bar{j}}(p) \\
=-\partial_{\bar{l}} \partial_{i} \ln (\operatorname{det}(h))(p) . \tag{2.5}
\end{gather*}
$$

The above calculation shows that $\operatorname{Ricc}(J Y, J Z)=\operatorname{Ricc}(Y, Z)$. Akin to $\omega$, we define a Ricci form on a Kähler manifold as $\operatorname{Ricc}(J X, Y)$. It is clear that this form is a $(1,1)$-form that is real (we will abuse notation and call this Ricci form also as the Ricci curvature sometimes). So the Ricci form is $\sqrt{-1} \bar{\partial} \partial \ln (\operatorname{det}(h))$. In other words it is simply $2 \pi c_{1}\left(K_{M}^{*}\right)$.

Let us now connect the Ricci tensor in these complex coordinates (that is acting in a Hermitian manner on $T^{1,0} M$ ) to real coordinates. In almost the same way as $h$ is related to $g$, the tensor $T(u, v)=\operatorname{Re}(\operatorname{Ricc}(\operatorname{Lu}, \overline{\operatorname{Lo}}))=\operatorname{Re}\left(\operatorname{Ricc}\left(\frac{u-\sqrt{-1} j u}{2}, \frac{v+\sqrt{-1} j v}{2}\right)\right)$ is equal to $\frac{1}{2} \operatorname{Ricc}(u, v)$. So the isomorphism gives a slightly different (by a factor of 2) Ricci tensor than the one we use in Riemannian geometry. The scalar curvature in Riemannian geometry is defined as the trace of the Ricci tensor.

Exercise : Prove that the scalar curvature in the complex setting above differs from the usual one by a factor of 4 .

The Riemannian sectional curvatures are $g(R(u, v) v, u)$. In the Kähler case, calculations above show that $R(x, y, z, w)=R(x, y, J z, J w)$. For a unit vector $x$, we define the holomorphic sectional curvature as $H(x)=R(x, J x, J x, x)$ and for two orthonormal unit vectors, $x, y$, the bisectional curvature is $R(x, J x, J y, y)$.

Exercise : Show that the bisectional curvature is $R(x, J x, J y, y)=R(x, y, y, x)+R(x, J y, J y, x)$ (thus justifying its name).

Note that $R_{j-\overline{i i}}^{-j--}=R\left(\partial_{i}, \partial_{\bar{i}}, \partial_{j}, \partial_{\bar{j}}\right)$ is the following using the usual isomorphism between $T^{1,0}$ and $T M$ (and symmetries of the Riemann tensor).

$$
\begin{equation*}
R_{j-i \bar{i}}^{-j--}=\frac{1}{4} R\left(\partial_{x^{i}} J \partial_{x^{i}}, J \partial_{x^{j}}, \partial_{x^{j}}\right) . \tag{2.6}
\end{equation*}
$$

Exercise : Show that all holomorphic sectional curvatures of $\mathbb{C}^{n}, \mathbb{C P}^{n}$, and $\mathbb{D}$ are 0 , positive constants, and negative constants respectively.

In the Riemannian case, when all the sectional curvatures are constant, the manifold is isometric to a quotient of Euclidean space, the Sphere, or Hyperbolic space. Akin to that, if all the holomorphic sectional curvatures are constant, the manifold is biholomorphically isometric to a quotient of $\mathbb{C}^{n}$, $\mathbb{C P}^{n}$, or $\mathbb{D}$.

