## MA 200 - Lecture 23

## 1 Recap

1. Orientations for 1-manifolds through unit tangent vector fields.
2. Orientations for hypersurfaces through unit normal vector fields.
3. Restrictions of orientations from $M$ to $\partial M$ produce orientations for $\partial M$. A standard orientation for $\partial M$ is the restriction if $\operatorname{dim}(M)$ is even, and the opposite of the restriction when $\operatorname{dim}(M)$ is odd.

## 2 Differential forms, wedge products, and form-fields in

 $\mathbb{R}^{n}$We want to generalise the notion of a cross product (because that will also help us generalise the notion of curl). Naively, if $a, b \in \mathbb{R}^{4}$, then $a \times b$ ought to have components $a_{i} b_{j}-a_{j} b_{i}$ where $i \neq j$. Even if we demand $i<j$ (because anyway, $a_{i} b_{j}-a_{j} b_{i}=-\left(a_{j} b_{i}-a_{i} b_{j}\right)$, i.e., $\left.a \times b=-b \times a\right)$, the number of components is 6 . So $a \times b$ cannot be a vector in $\mathbb{R}^{4}$ ! (On the other hand, $a \times b \times c$ in this naive prescription will have only 4 components!) So we need to extend our definitions from vectors to other beasts. Whatever this naive $a \times b$ is, each component is certainly multilinear in $a, b$ and antisymmetric. That motivates the following definitions.
Let $V$ be a f.d real vector space. We will almost always consider $V=\mathbb{R}^{k}$ or $T_{p} M$ for some manifold-with-boundary. We know what a multilinear map $T: V \times V \times \ldots \rightarrow \mathbb{R}$ is (what is it?) Multilinear maps from $V^{m}$ to $\mathbb{R}$ are also called $m$-tensors. An example is $T(\vec{v}, \vec{w})=\operatorname{det}(\vec{v} \vec{w})$ where $\vec{v}, \vec{w} \in \mathbb{R}^{2}$. Another is $p(\vec{v}, \vec{w})=v_{1} w_{1}+v_{2} w_{2}=\langle v, w\rangle$. The set of $m$-tensors forms a vector space. We can find a nice basis for this vector space: Given an $m$-multiindex $I=\left(i_{1}, \ldots, i_{m}\right)$ and a basis $e_{1}, \ldots, e_{n}$, the tensors $\phi_{I}\left(v_{1}, \ldots, v_{m}\right)=\left(v_{1}\right)_{i_{1}}\left(v_{2}\right)_{i_{2}} \ldots$ form a basis (so the dimension is $n^{m}$ where $n=\operatorname{dim}(V)$ ). By the way, for a single index, the 1-tensors $\phi_{i}$ are called the dual basis of $e_{j}$. 1-tensors are also called linear functionals, or as we shall call them later, 1 -forms. (By the 0 -forms are simply real numbers.)
A neat way of phrasing this statement is through the definition of the tensor product of two tensors: Suppose $S, T$ are $k$ and $l$ tensors respectively. The map $S \otimes T$ defined as $S \otimes T\left(v_{1}, \ldots, v_{k}, w_{1}, \ldots, w_{l}\right)=S(v) T(w)$ is a $(k+l)$-tensor called the tensor product of $S$ and $T$. Note that $\phi_{I}=\phi_{i_{1}} \otimes \phi_{i_{2}} \ldots$. This tensor product obeys some standard properties, namely, associativity and linearity.

Given a linear map $T: V \rightarrow W$, there is a 'dual' linear map denoted as $T^{*}$ from $k$-tensors on $W$ to those on $V: T^{*} L\left(v_{1}, \ldots, v_{k}\right)=L\left(T v_{1}, \ldots, T v_{k}\right)$. This dual map obeys some (easy to verify) properties: $T^{*}$ is linear on the space of $k$-tensors on $W$, $T^{*}(f \otimes g)=T^{*} f \otimes T^{*} g$ and $(S \circ T)^{*}=T^{*} \circ S^{*}$.
What we want to study are tensors that are antisymmetric or alternating (akin to our example of determinants, or cross products):
Def: A $k$-tensor is said to be symmetric if $f\left(v_{1}, \ldots, v_{k}\right)=f\left(v_{1}, \ldots, v_{i+1}, v_{i}, \ldots\right)$, i.e., interchanging adjacent arguments doesn't change the value. By a finite number of adjacent interchanges, one can see that interchanging any two arguments keeps the value invariant. Since every permutation $\sigma \in S_{k}$ is a product of transpositions (hopefully you did this in UM 205), the value is invariant under any permutation of the arguments. A $k$-tensor is said to be antisymmetric or alternating $f\left(v_{1}, \ldots, v_{k}\right)=-f\left(v_{1}, \ldots, v_{i+1}, v_{i}, \ldots\right)$. Such tensors are also called $k$-forms and they form a vector space that is denoted as $\Lambda^{k} V$. (For $k=1$, every tensor is trivially alternating.) They are also called $k$-forms. Clearly, if $v_{i}=v_{j}$ for some $i \neq j$, the value of $f$ is zero on such a tuple (if $|j-i|=1$, this is trivial. Now induct on $|j-i|$ ). Moreover, if a tensor is such that it is zero whenever adjacent arguments coincide, then it is alternating: $0=f\left(v_{1}, \ldots, v_{i}+v_{i+1}, v_{i}+v_{i+1}, v_{i+2}, \ldots\right)$. Now use multi-linearity and the vanishing property again to be done.
Inner products are examples of symmetric 2-tensors and the determinant an example of $\Lambda^{n}\left(\mathbb{R}^{n}\right)$. In fact, we can produce alternating 3 -tensors in $\mathbb{R}^{4}$ using the determinant as $(u, v, w) \rightarrow \operatorname{det}\left(e_{1} u v w\right)$ and so on. So what is the dimension of $\Lambda^{k} V$ ? Suppose $e_{1}, \ldots, e_{n}$ is a basis of $V$. Define the following alternating $k$-tensors for a given $k$-multiindex $I=\left(i_{1}, \ldots, i_{k}\right)$ :

$$
\begin{equation*}
\epsilon_{I}\left(v_{1}, \ldots, v_{k}\right)=\operatorname{det}\left(A_{I}\right) \text { where }\left(A_{I}\right)_{j l}=\left(v_{l}\right)_{i_{j}} . \tag{1}
\end{equation*}
$$

Here are examples:

1. When $k=1, \epsilon_{i}$ form the dual basis for $V^{*}=\Lambda^{1} V$.
2. Suppose $k=2$ and $I=(2,3), v_{1}=(a, b, c), v_{2}=(\alpha, \beta, \gamma)$ then $\epsilon_{I}\left(v_{1}, v_{2}\right)=$ $\operatorname{det}\left(\begin{array}{ll}b & \beta \\ c & \gamma\end{array}\right)=b \gamma-c \beta$.
3. Suppose $i_{1}=i_{3}$. Then $\epsilon_{I}=0$ (why?) More generally, when two indices coincide, $\epsilon_{I}=0$.
4. Recall that the sign of a permutation $\sigma \in S_{k}$ is simply the determinant of the corresponding permutation matrix. Now we define $\sigma(I)=\left(i_{\sigma(1)}, \ldots\right.$ ). (For instance, if $k=3$, and $\sigma=(12)$, then $\sigma(2,3,1)=(3,2,1)$. Then $\epsilon_{\sigma(I)}=\operatorname{sgn}(\sigma) \epsilon_{I}$ by the properties of determinants.

Now we consider the multiindices $I$ such that $i_{1}<i_{2}<\ldots<i_{k}$. There are $\binom{n}{k}$ of these indices. We claim that such $\epsilon_{I}$ form a basis for $\Lambda^{k}(V)$. Indeed,

1. These $\epsilon_{I}$ are linearly independent: Suppose $\sum_{i_{1}<\ldots<i_{k}} c_{I} \epsilon_{I}=0$. Then if $j_{1}<j_{2}<$ $\ldots<j_{k}, 0=\sum c_{I} \epsilon_{I}\left(e_{j_{1}}, \ldots, e_{j_{k}}\right)=c_{J}$ (why?)
2. They span $\Lambda^{k}(V)$ : Let $\omega \in \Lambda^{k}(V)$. Then consider $\tilde{\omega}=\sum \omega\left(e_{i_{1}}, \ldots, e_{i_{k}}\right) \epsilon_{I}$. Note that to prove that $\tilde{\omega}=\omega$, it is enough to show that they are equal on $\left(e_{j_{1}}, \ldots, e_{j_{k}}\right)$ for all increasing multiindices $J$ (why? Because of multilinearity, and the facts that if two indices coincide, we get zero and if we change the ordering, we pick up the sign of the permutation). Now $\tilde{\omega}\left(e_{J}\right)=\sum \omega\left(e_{I}\right) \epsilon_{I}\left(e_{J}\right)=\omega\left(e_{J}\right)$ (where we abuse notation by denoting a tuple $\left(w_{i_{1}}, \ldots\right)$ by $\left.w_{I}\right)$.

Note that when $k=n, \Lambda^{n}$ has dimension 1 and is generated by $\epsilon_{12 \ldots n}$. These forms are also called 'top forms'. When $k=0$, again the dimension is 1 . For $n=3, \Lambda^{1}$ and $\Lambda^{2}$ have exactly the same dimension equal to 3 (which is also the dimension of $V$ !) So we can identify a vector $v \in \mathbb{R}^{3}$ with a 1-form $\omega=v_{1} \epsilon_{1}+v_{2} \epsilon_{2}+v_{3} \epsilon_{3}$ and with a 2-form $v_{1} \epsilon_{23}+v_{2} \epsilon_{31}+v_{3} \epsilon_{12}$ (why this weird identification? Because we want to think of $\epsilon_{12}$ as $\hat{i} \times \hat{j}=\hat{k})$.
We are now in a position to define the generalisation of the cross product. Instead of defining it for vectors, we define the wedge product $\wedge$ of forms. We want to try the following naive definition: Let $\omega \in \Lambda^{k}(V)$ and $\eta \in \Lambda^{l}(V)$. Then $\omega \wedge \eta \in \Lambda^{k+l}(V)$ is defined as $\omega \wedge \eta=\sum_{i_{1}<\ldots<i_{k}, j_{1}<\ldots<j_{l}} \omega_{I} \eta_{J} \epsilon_{I J}$, i.e., we define $\epsilon_{I} \wedge \epsilon_{J}$ as $\epsilon_{I J}$ and extend this definition linearly.

