

• HESSIAN OF THE ENERGY.

• Recall: 1) Given an admissible curve $\gamma: [a, b] \rightarrow M \in \Omega(p, q)$, the length & energy are defined as.

$$L(\gamma) = \sum_{k=1}^n \int_{a_{k-1}}^{a_k} |\dot{\gamma}(t)| dt$$

$$E(\gamma) = \frac{1}{2} \sum_{k=1}^n \int_{a_{k-1}}^{a_k} |\dot{\gamma}(t)|^2 dt$$

where $\gamma|_{[a_{k-1}, a_k]}$ is smooth.

2) If $\Gamma: (-\varepsilon, \varepsilon) \times [a, b] \rightarrow M$ is an admissible variation, then for $\gamma_s = \Gamma(s, \cdot)$ & $V_s = \partial_s \Gamma(s, \cdot)$ we have

$$\begin{aligned} \frac{d}{ds} E(\gamma_s) &= - \int_a^b \langle \partial_s \Gamma, D_t \dot{\gamma}_s \rangle + \langle V_s(b), \dot{\gamma}_s(b-) \rangle \\ &\quad - \langle V_s(a), \dot{\gamma}_s(a+) \rangle - \sum_{k=1}^m \langle V_s(a_k), \Delta_k \dot{\gamma}_s \rangle \end{aligned}$$

where $\Delta_k W := W(a_{k+}) - W(a_{k-})$.

Defⁿ: An admissible 2-parameter variation of γ is a map $\Gamma: (-\varepsilon, \varepsilon)^2 \times [a, b] \rightarrow M$ s.t.

1) \exists a sub-division $a = a_0 < a_1 < \dots < a_m = b$ s.t. Γ is smooth on $(-\varepsilon, \varepsilon)^2 \times [a_{k-1}, a_k] \forall k$.

2) $\gamma_0(t) := \Gamma(0, 0, t) = \gamma(t)$.

We set $\gamma_{s_1, s_2} := \Gamma(s_1, s_2, \cdot)$.

We denote $W_i(t) = \partial_{s_i} \Gamma(0, 0, t)$ for $i = 1, 2$. (2)

We say Γ is proper if $\Gamma(s_1, s_2, a) = \gamma(a)$ &
 $\Gamma(s_1, s_2, b) = \gamma(b) \quad \forall (s_1, s_2) \in (-\varepsilon, \varepsilon)^2$. (\Leftarrow)

In particular $W_i(a) = 0, W_i(b) = 0$.

Ex: If W_1, W_2 are v.f. along γ , then

$$\Gamma(s_1, s_2, t) := \exp_{\gamma(t)}(s_1 W_1(t) + s_2 W_2(t))$$

is an example of such a variation. If $W_i(a) = 0, W_i(b) = 0$, then Γ is proper.

Th^m 19.1 (Second variation formula). Let Γ be an admissible 2-parameter variation of a geodesic γ w/ variation vector fields $W_i(t); i=1, 2$.

Then

$$\begin{aligned} \frac{\partial^2 \mathcal{E}}{\partial s_1 \partial s_2} \Big|_{(s_1, s_2) = (0, 0)} \mathcal{E}(\Gamma_{s_1, s_2}) &= \int_a^b \left(\langle D_t W_1, D_t W_2 \rangle - \text{Rm}(W_1, \dot{\gamma}, \dot{\gamma}, W_2) \right) dt \\ &= - \sum_{k=1}^{n-1} \langle \Delta_k D_t W_1, W_2(a_k) \rangle \\ &\quad - \int_a^b \left\langle \frac{D^2 W_1}{dt^2} + R(W_1, \dot{\gamma}) \dot{\gamma}, W_2 \right\rangle dt \end{aligned}$$

(3)

Pf: let $E_k(s_1, s_2) = \mathcal{E}(r_{s_1, s_2} | [a_{k-1}, a_k])$.

Then

$$\begin{aligned} \frac{\partial E_k}{\partial s_1} &= \int_{a_{k-1}}^{a_k} \left\langle \frac{\mathcal{D}}{\partial s_1} \partial_t \Gamma, \partial_t \Gamma \right\rangle \\ &= \int_{a_{k-1}}^{a_k} \left\langle \frac{\mathcal{D}}{\partial t} \partial_{s_1} \Gamma, \partial_t \Gamma \right\rangle \quad (\text{Symmetry lemma}). \end{aligned}$$

Differentiating again

$$\frac{\partial^2 E_k}{\partial s_1 \partial s_2} = \int_{a_{k-1}}^{a_k} \left\langle \mathcal{D}_{s_2} \mathcal{D}_t \partial_{s_1} \Gamma, \partial_t \Gamma \right\rangle + \int_{a_{k-1}}^{a_k} \left\langle \mathcal{D}_t \partial_{s_1} \Gamma, \mathcal{D}_t \partial_{s_1} \Gamma \right\rangle.$$

For the first term, if we put $\partial_{s_i} \Gamma = S_i$, $\partial_t \Gamma = T$,

$$\mathcal{D}_{s_2} \mathcal{D}_t \partial_{s_1} \Gamma = \mathcal{D}_t \mathcal{D}_{s_2} \partial_{s_1} \Gamma + R(S_2, T) S_1.$$

$$\text{So, } \int_{a_{k-1}}^{a_k} \left\langle \mathcal{D}_{s_2} \mathcal{D}_t \partial_{s_1} \Gamma, \partial_t \Gamma \right\rangle = \int_{a_{k-1}}^{a_k} \left\langle \mathcal{D}_t \mathcal{D}_{s_2} S_1, T \right\rangle - \int R_m(S_2, T, T, S_1).$$

$$\text{Since } \mathcal{D}_t T \Big|_{(s_1, s_2)} = \frac{\partial}{\partial t} \left\langle \mathcal{D}_{s_2} S_1, T \right\rangle - \int_{a_{k-1}}^{a_k} \left\langle \mathcal{D}_{s_2} S_1, \mathcal{D}_t T \right\rangle$$

$$= - \int R_m(S_2, T, T, S_1)$$

$$\begin{aligned} &= \left\langle \mathcal{D}_{s_2} S_1(a_k^-), T(a_k^-) \right\rangle - \left\langle \mathcal{D}_{s_2} S_1(a_{k-1}^+), T(a_{k-1}^+) \right\rangle \\ &\quad - \int_{a_{k-1}}^{a_k} \left\langle \mathcal{D}_{s_2} S_1, \mathcal{D}_t T \right\rangle - \int_{a_{k-1}}^{a_k} R_m(S_2, T, T, S_1) \end{aligned}$$

Plugging in $(\delta_1, \delta_2) = 0$, since r is a geodesic. (4)

$$D_t T|_{(0,0)} = D_t \dot{r} = 0. \text{ Also, } S_1(0,0) = W_1, S_2(0,0) = W_2.$$

So,

$$\left. \frac{\partial^2 \mathcal{E}}{\partial \delta_1 \partial \delta_2} \right|_{(0,0)} = \Delta + \int_a^b \langle D_t W_1, D_t W_2 \rangle - R_m(W_1, \dot{r}, \dot{r}, W_2)$$

Claim 1: $\Delta \equiv 0$

Pf: Note $\Delta = \langle D_{\delta_2} S_1(b), T(b-) \rangle - \langle D_{\delta_2} S_1(a+), T(a+) \rangle$
 $+ \sum_{k=1}^{n-1} \langle D_{\delta_2} S_1(a_{k+}), T(a_{k+}) \rangle - \langle D_{\delta_2} S_1(a_{k-}), T(a_{k-}) \rangle$

Since $S_1(\delta_1, \delta_2, a) \equiv 0$ & $S_1(\delta_1, \delta_2, b) \equiv 0$, 1st 2 terms are zero. Also at $t = a_k$, $D_{\delta_2} S_1 = D_{\delta_2} \partial_{\delta_1} \Gamma$ depends only on values of Γ on $t = a_k$. So $D_{\delta_2} S_1$ is cont. $\forall (s_1, s_2, t)$. Also at $(\delta_1, \delta_2) = (0,0)$ $T = \dot{r}$. Since r is an unbroken geodesic, $\dot{r}(a_{k+}) = \dot{r}(a_{k-})$. So, $\Delta \equiv 0$.

For the 2nd formula, we have

Claim 2: $\int_a^b \langle D_t W_1, D_t W_2 \rangle = - \int \langle \frac{D^2 W_1}{dt^2}, W_2 \rangle$
 $- \sum_{k=1}^n \langle \Delta_k D_t W_1, W_2(a_k) \rangle$

(5)

Pf:
$$\int_{a_{k-1}}^{a_k} \langle \mathcal{D}_t W_1, \mathcal{D}_t W_2 \rangle = \int_{a_{k-1}}^{a_k} \frac{\partial}{\partial t} \langle \mathcal{D}_t W_1, W_2 \rangle - \int_{a_{k-1}}^{a_k} \langle \frac{\mathcal{D}^2 W_1}{dt^2}, W_2 \rangle$$

$$= \langle \mathcal{D}_t W_1(a_k^-), W_2(a_k^-) \rangle - \langle \mathcal{D}_t W_1(a_{k+1}^+), W_2(a_{k+1}^+) \rangle - \int_{a_{k-1}}^{a_k} \langle \frac{\mathcal{D}^2 W_1}{dt^2}, W_2 \rangle$$

Again $\partial_{s_2} \Gamma$ is cont. at $t = a_k$. So $W_2(a_{k+1}^+) = W_2(a_k^-)$.
 So adding up, we get the desired result.

Prop 19.2 If γ is a unit speed geodesic & Γ is a 2-parameter admissible ^{prop.} variation, then

$$\frac{\partial^2 L(\gamma_{s_1, s_2})}{\partial s_1 \partial s_2} \Big|_{(0,0)} = \int_a^b \langle \mathcal{D}_t W_1, \mathcal{D}_t W_2 \rangle - \langle \mathcal{D}_t W_1, \dot{\gamma} \rangle \langle \mathcal{D}_t W_2, \dot{\gamma} \rangle - Rm(W_1, \dot{\gamma}, \dot{\gamma}, W_2)$$

$$= \int_a^b \langle \mathcal{D}_t W_1^\perp, \mathcal{D}_t W_2^\perp \rangle - Rm(W_1^\perp, \dot{\gamma}, \dot{\gamma}, W_2^\perp)$$

where W_i^\perp is the component of W_i in $N\gamma(t)$

Pf: We only prove the final equality. Writing $W = W^T + W^\perp$ we have $W^T = \langle W, \dot{\gamma} \rangle \dot{\gamma}$.

So $\mathcal{D}_t W^T = \langle \mathcal{D}_t W, \dot{\gamma} \rangle \dot{\gamma} = (\mathcal{D}_t W)^T$.

$\mathcal{D}_t W^\perp = (\mathcal{D}_t W)^\perp$.

$$\text{So } \langle D_t W_1, D_t W_2 \rangle = \langle (D_t W_1)^T, (D_t W_2)^T \rangle \quad (6)$$

$$+ \langle D_t W_1^\perp, D_t W_2^\perp \rangle.$$

$$= \langle D_t W_1, \dot{\gamma} \rangle \langle D_t W_2, \dot{\gamma} \rangle$$

$$+ \langle D_t W_1^\perp, D_t W_2^\perp \rangle.$$

Also since $Rm(\dot{\gamma}, \dot{\gamma}, \cdot, \cdot) = Rm(\cdot, \cdot, \dot{\gamma}, \dot{\gamma})$, and

$\langle W_i^T, \dot{\gamma} \rangle = 0$, we have $Rm(W_i, \dot{\gamma}, \dot{\gamma}, W_2) = Rm(W_i^\perp, \dot{\gamma}, \dot{\gamma}, W_2^\perp)$

Defⁿ: We define the "tangent sp" to $\Omega(p, q)$ at $\gamma \in \Omega(p, q)$ by

$$T_\gamma \Omega(p, q) := \{W \in \mathcal{F}(\gamma) \mid W(a) = 0, W(b) = 0\}$$

2) The Hessian of $E: \Omega(p, q) \rightarrow \mathbb{R}$ at a geodesic γ is defined by $\text{Hess}_\gamma E: T_\gamma \Omega \times T_\gamma \Omega \rightarrow \mathbb{R}$.

$$\text{Hess}_\gamma E(W_1, W_2) := \int_a^b \langle D_t W_1, D_t W_2 \rangle - Rm(W_1, \dot{\gamma}, \dot{\gamma}, W_2)$$

||

$I_\gamma(W_1, W_2) \leftarrow$ index form.

Prop 19.3 I_γ is a symmetric, bilinear form on $T_\gamma \Omega(p, q)$

Cor 19.4: Let γ be a minimal geodesic between p and q . Then $I_\gamma(W, W) \geq 0 \quad \forall W \in T_\gamma \Omega$.

Pf: Let Γ be any variation w/ $\partial_s \Gamma(0, t) = W(t)$
(e.g. $\Gamma(s, t) = \exp_{\gamma(t)} s W(t)$).

Let $r_s(t) = \Gamma(s, \cdot)$. Then Cauchy Schwartz & r being minimal \Rightarrow

$$\mathcal{E}(r_s) \geq \frac{1}{2(b-a)} L(r_s)^2 \geq \frac{1}{2(b-a)} L(r)^2 = \mathcal{E}(r).$$

So $\mathcal{E}(\Gamma(s, \cdot))$ has a minimum at $s=0$. So

$$0 \leq \left. \frac{d^2}{ds^2} \right|_{s=0} \mathcal{E}(r_s) = I_r(W, W).$$

JACOBI EQUATION:

Defⁿ: The null space of a geodesic r is defined

by $\text{Null}_r := \{W \in T_r \Omega \mid I_r(W, W') = 0 \forall W'\}$

Prop 19.5: $J \in \text{Null}_r$ if and only if J is smooth

& satisfies

$$\frac{D^2 J}{dt^2} + R_{\alpha\beta}(J, \dot{r}) \dot{r} = 0 \quad (*)$$

Pf: \Leftarrow by the second variation formula, since J is smooth (& hence $\Delta_{\dot{r}} D_t J = 0$) we have $J \in \text{Null}_r$.

\Rightarrow : Conversely sp $J \in \text{Null}_r$, then choose a partition $a = a_0 < \dots < a_m = b$ s.t. $J|_{[a_{k-1}, a_k]}$ is

C^∞ . Let $\varphi: [a, b] \rightarrow [0, 1]$ smooth s.t.

(1) $\varphi(a_k) = 0 \forall k$

(2) $\varphi > 0$ otherwise.

and let

(8)

$$W := \begin{cases} \varphi(t) \left(\frac{D^2 J}{dt^2} + R(J, \dot{r}) \dot{r} \right), & t \neq a_k \\ 0 & \text{otherwise} \end{cases}$$

Then $0 = -I_r(J, W) = \int_a^b \varphi(t) \left[\frac{D^2 J}{dt^2} + R(J, \dot{r}) \dot{r} \right]$

So $\frac{D^2 J}{dt^2} + R(J, \dot{r}) \dot{r} = 0$ for $t \neq a_k$.

Claim: J is smooth.

Pf: Now let $W' \in T_r \Omega$ s.t.

$$W'(a_k) = \Delta_k \frac{DJ}{dt}$$

Then

$$0 = -I_r(J, W') = \sum_k \left\| \Delta_k \frac{DJ}{dt} \right\|^2$$

So $\Delta_k \frac{DJ}{dt} = 0 \forall k$ i.e. $\frac{DJ}{dt}$ is cont at a_k .

From Th^m 20.1 (to be proved next time) any solution to $*$ on (a_k, a_{k+1}) is completely determined by $J(a_k)$ & $DJ/dt(a_k)$. Hence J is smooth at a_k .

D.f'' A

Defⁿ: A vector field J along γ is called a Jacobi field if (9)

$$\frac{D^2 J}{dt^2} + R_{\bullet}(J, \dot{\gamma}) \dot{\gamma} = 0.$$

Rk: We do not impose the restriction that $J(a) = 0$ & $J(b) = 0$. So $J \notin T_r \Omega$. Prop can be reformulated as:

$$J \in \text{Null}_r \iff J \text{ is a Jacobi field w/ } J(a) = 0 \\ J(b) = 0.$$

Let \mathcal{H} be a Hilbert space and T a self-adjoint operator on \mathcal{H} .
Then $\mathcal{H} = \mathcal{H}_+ \oplus \mathcal{H}_0 \oplus \mathcal{H}_-$ where $\mathcal{H}_\pm = \ker(T \mp |T|)$ and $\mathcal{H}_0 = \ker |T|$.

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